## Part 2: Lensing by distributed masses

- Theory
- Cluster mass
- Subhalos
- Applications:
- Search for subhalos
- Transients: SN Refsdel and Icarus
- Cosmography



## Part 2: Lensing by distributed masses - Cluster

- Consider an extended, axially-symmetric lens. Deflections are radial \& the mass $M(<b)$ matters:

$$
\hat{\alpha}=\frac{4 G M(<b)}{c^{2} b}
$$

- Recall the lens equation: $\theta-\beta=\frac{d_{L S}}{d_{S}} \hat{\alpha}$. The deflection is no longer simple. On substituting:

$$
\theta-\beta=\frac{d_{L S}}{d_{S}} \frac{4 G M}{c^{2} b}
$$



- Let us say this lens has constant surface mass density, $\Sigma$, such that $M=\Sigma \pi b^{2}$. On making this substitution, recalling that $b=d_{L} \theta$, and solving for $\beta$ :

$$
\beta=\theta\left(1-\frac{4 \pi G \Sigma}{c^{2}} \frac{d_{L} d_{L S}}{d_{S}}\right)
$$

- Now, $\beta \sim \theta$. If we set $\Sigma_{c r}=\frac{c^{2}}{4 \pi G} \frac{d_{S} d_{L}}{d_{L S}}$, then $\beta=0$ for all $\theta$. When $\Sigma=\Sigma_{\mathrm{cr}}$, we get giant arcs


## Part 2: Lensing by distributed masses - Cluster

-For the canonical singular isothermal sphere:

$$
\rho(r)=\frac{v_{c}^{2}}{4 \pi G r^{2}}, \text { which has projected surface mass density: } \Sigma(r)=\frac{v_{c}^{2}}{4 G r}
$$

- Then the mass interior is: $M=\int_{0}^{R} \Sigma(r) 2 \pi r d r=v_{c}^{2} \pi R / 2 G$
- So the deflection angle is: $\hat{\alpha}=\frac{4 G M}{c^{2} R}=\frac{2 \pi v_{c}^{2}}{c^{2}}<==$ independent of radius!
- The lens equation is: $\theta-\beta=\frac{d_{L S}}{d_{S}} \frac{2 \pi v_{c}^{2}}{c^{2}}= \pm \theta_{E} . \beta=0$ gives the Einstein ring. Otherwise, again we get $+/$ images
- There is an $\mathrm{I}_{+}$at $\beta>\theta_{\mathrm{E}}$ and I - at $\beta<\theta_{\mathrm{E}}$. $\mathrm{I}_{+}$is brighter than $I_{-}$, and $I_{\text {- is parity flipped. Note }}$ for $\theta>2 \theta_{\mathrm{E}}$ there is only one image



## Part 2: Lensing by distributed masses - Cluster

- Here is an example of the imaging of an extended source by a nonsingular circularly-symmetric lens.
- A source near to the center produces two tangentially-stretched images, and a source on the caustic produces one radial arc and one tangential arc.

- real mass distributions are not circularly-symmetric. For this elliptical lens, there can be 4 or even more images. Here source plane has regions where odd-numbered images are created, yet one image is demagnified



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## Part 2: Lensing by distributed masses - Cluster

- Arc locations give projected cluster mass inside circle traced by the arc
- For a circularly-symmetric lens, $\Sigma$ is approximately equal to $\Sigma_{\text {cr }}$
-The radius gives an estimate of Einstein radius of the cluster:

$$
M(\theta)=\Sigma_{c r} \pi\left(d_{L} \theta\right)^{2} \approx 1.1 \times 10^{14} M_{\odot}\left(\frac{\theta}{30^{\prime \prime}}\right)^{2}\left(\frac{D}{1 G p c}\right)
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The "Sextet Arcs"
A1689 (z=0.18);
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G165 with critical curve and arclet families overlaid: this image was reduced, and the lens model fit made, by Massimo!

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## Part 2: Lensing by distributed masses - Subhalo

- Dark matter is an outstanding problem: where/what is it? ${ }^{\text {I }}$ solated masses with $M=10^{7}-10^{8}$ Msun have $\Sigma \ll \Sigma_{\text {cr }}$
- In a cluster, we are given the position of a lensed image $\vec{x}$ whose coordinates are altered from its source coordinates $\vec{y}: \vec{y}=A \vec{x}$
- How do we envision this mapping, as $\alpha$ depends on $\theta$ nonlinearly? We zoom to vicinity of point x . The best-fit linear map is the Jacobian
- We choose coordinates to be eigenvectors, lens plane which means also that A is diagonal. We impose the condition that on the critical $\operatorname{curve} \operatorname{det}(\mathrm{A})=$ 0 . We approximate A:

$$
A=\left|\begin{array}{cc}
\vec{d} \cdot \vec{x} & 0 \\
0 & 2\left(1-\kappa_{0}\right)
\end{array}\right|
$$



12 (Venumadhav et al. 2018; Dai et al. 2018)

## Part 2: Lensing by distributed masses - Subhalo

- Recall the cluster lens equation: $\vec{y}=\vec{x}-\vec{\alpha}_{B}(\vec{x})$ Substituting into A:

$$
A=\left[\frac{\partial \vec{y}}{\partial \vec{x}}\right]=\left[\frac{\partial\left(\vec{x}-\vec{\alpha}_{B}(\vec{x})\right)}{\partial \vec{x}}\right]=1-\left[\frac{\partial \vec{\alpha}_{B}(\vec{x})}{\partial \vec{x}}\right]
$$

- Let's introduce a subhalo which causes a deflection angle $\vec{\alpha}_{s h}$ which shifts the image to a new position of $\vec{x}^{\prime}=\vec{x}+\Delta x$. With a subhalo:

$$
\vec{y}=\vec{x}^{\prime}-\vec{\alpha}_{B}\left(\vec{x}^{\prime}\right)-\vec{\alpha}_{s h}(\vec{x})
$$

- In each case, the y -vector is is the same. On setting the two lens equations equal to each other and Taylor-expanding the term $\vec{\alpha}_{B}(\vec{x}+\Delta x)$

$$
\begin{gathered}
\vec{x}^{\prime}-\vec{x}-\left(\vec{\alpha}_{B}\left(\vec{x}^{\prime}\right)-\vec{\alpha}_{B}(\vec{x})\right)=\vec{\alpha}_{s h} \\
\Delta x\left(1-\left[\frac{\partial \overrightarrow{\alpha_{B}}}{\partial \vec{x}}\right]\right)=\vec{\alpha}_{s h} \\
\Delta x=A^{-1}\left(\vec{x}^{\prime}\right) \vec{\alpha}_{s h}
\end{gathered}
$$

## Part 2: Lensing by distributed masses - Subhalo

- The shift in position as a result of the subhalo is: $\Delta x=A^{-1}\left(\vec{x}^{\prime}\right) \vec{\alpha}_{s h}$
- Recall also the Jacobian: $A^{-1}=\left|\begin{array}{cc}\frac{1}{\vec{d} \cdot \vec{x}} & 0 \\ 0 & 1 / c\end{array}\right|$, where $c=2\left(1-\kappa_{0}\right)$
- The magnification factor: $\mu=1 / \operatorname{det}(A)=\frac{1}{c \cdot \vec{x} \cdot \vec{d}}$
- The shifts in the x 1 and x 2 directions are:

$$
\begin{aligned}
\Delta x_{1} & =\frac{1}{\vec{x} \cdot \vec{d}} \cdot \alpha_{1, s h}\left(x^{\prime}\right)=C \mu \alpha_{1, s h}\left(x^{\prime}\right) \\
\Delta x_{2} & =\frac{1}{C} \alpha_{2, s h}\left(x^{\prime}\right)
\end{aligned}
$$

critical curve

- The image is stretched in $x_{1}$ direction, and is a double. Compact source images found close to main critical curve identifies subhalos!


## Part 2: Lensing by distributed masses - Subhalo

- An image pair enhanced by subhalo will be stretched along $x_{1}$ direction, bridging the critical curve.
- Source must be compact, as a galaxy is too extended.
- Images must be local, $\sim 0.1$ arcsec from critical curve!

Arc 1a in G165

(Frye et al. 2018)

- The "dragon" in A370 is good for its extensive "tail" crossing critical curve - "Arc 1a" in G165 is an SMG, and may already have blue compact star forming knots. Is this an image pair?

The Dragon in A370

(Dai et al. 2018)

MACS J1149 SN Refsidel, $z=1.52$

Galaxy bost appears in 4 different locations!

SN Refsdel appears in 4 locations in the brigbtest image, $m_{A B}=25.5 \mathrm{mag}$
(Kelly et al. 2016, 2018)

## Applications: transients



- Why search for strongly-lensed SNe?
- standard candles offer independent test of image magnification
- multiply-imaged SNe count as arclet families for cosmography
- for ideal geometries, time delays $==>\mathrm{H}_{0}$

Caustic transit at $\mathrm{z}=1.53$

- Why search for caustic transit events?
- ultra-high magnification factors of 103-104
- sensitivity to microlensing effects
- route to find first light sources, although event rates depend on ICL, background population, and luck

(Kelly+18)


## Applications: cosmography

- Constraining Cosmological Parameters, $w_{x}$
- Multiple sets of arclet families sensitive to the geometry $==>$ ratio of angular diameter distances
- Take away: 2-3 arclet families in $\mathbf{1 0 0}$ clusters with redshifts and lens models yield constraints competitive with other methods, and samples different parameter space





(Gilmore \& Natarajan+09)


## Long Duration Transients with LSST: Getting Acquainted with our Dim Neighbors

$\downarrow$ Verification - of 20,000 annual events, 1000 s will be degenerate objects. LSST will help break mass/velocity degeneracy that turns measurement tross into measurements of mass
$\uparrow$ For $\alpha \sim 0.1$-few, LSST+GAIA+Pan-STARRS enable detection of BH-BH binaries at physical separations are extensions of the ones that will merge to produce LIGO events
\& LSST annual event rates ${ }^{1}: 14,000 \mathrm{M}$ dwarfs, 2,600 WDs,
$2 \theta_{E}=0.06^{\prime \prime}$
$\longleftrightarrow$
$\ddots$
$\vdots$
$\omega=0.2^{\prime \prime} y r^{-1}$
$\tau_{c}=100 d a y s$
(Stellar BH) 1,800 neutron stars, $120 \mathrm{BHs}, \mathrm{RR}$ Lyrae/other objects.
† On WFIRST: which are lessons learned from space-based studies?
${ }^{1}$ from Di Stefano 2008a, consistent with MACHO event rate extrapolation

